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PHOTOLUMINESCENCE OBSERVATIONS OF DONOR-ACCEPTOR EMISSION IN MILLED Cd<sub>0.3</sub>Zn<sub>0.7</sub>Se NANOCRYSTALS

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### ABSTRACT

Defects and impurities that arise during material synthesis are detrimental to device performance. Hence, an effective characterization tool is fundamental to effectively quantify and understand the origin of such defects. In this paper, we examine photoluminescence spectra (PL) from (650-900) nm of mechanically alloyed  $Cd_{0.3}Zn_{0.7}Se$  nanocrystals and observe three prominent emission peaks (1.4, 1.8, and 1.54) eV. The peaks were pinned at a fixed position despite a decrease in crystallites size which ought to increase the energy bandgap. The intensity of the peaks increases with milling time due to longer milling time which gives rise to a higher density of intrinsic defects. This was attributed to continuous deformations from the milling process. The formation energies of acceptor and donor defects remain constant to the conduction and valance band edge.

Keywords: Mechanical alloying, Photoluminescence spectra (PL), Donor-acceptor pairs (DAP), defects.

## INTRODUCTION

Alloyed CdZnSe nanocrystal has been employed for the fabrication of light-emitting devices superior with spectral properties. The bottle-neck surrounding its wider applications in areas such as absorber layers in solar cells, sensors, and biological labeling is the presence of non-radiative decay channels associated with surface state defects, intrinsic point defects, and other stoichiometric defects.

The presence of deep-level emission (DLE) in undoped CdZnSe involves a hole of the intrinsic point defects such as vacancies, interstitials, and antisites. It is well known that point defects affect the electrical and optical properties of CdZnSe. There have been limited reports on the DL emission in CdZnSe nanocrystals. (Bhattacharyya et al., 2009) observed a broad DL emission from 590-750 nm (2.10 to 1.65 eV) centered at 650 nm (1.839 eV) which was attributed to radiative recombination involving donor-acceptor pairs (DAP). Zhang et al. (Zhang & Hark, 2000), reported DL emission from (1.3 to 1.8 eV) in Zn<sub>x</sub>Cd<sub>1-x</sub>Se grown by organometallic-vapor-phase epitaxy (OMVPE). The DL emission originated from cation-vacancy related to stacking faults. In the same vein, Zn<sub>x</sub>Cd<sub>1-x</sub>Se grown by metal-organic chemical vapor deposition (MOCVD) on InP substrate, produced a broad DL emission observed from (1.35-1.9 eV). The origin of DLE was also attributed to cation-vacancy related complexes (Zhang et al., 1998). However, Burger et al. (Burger & Roth, 1984), reported DLE at 0.49 eV below the conduction band in Zn0.3Cd0.7Se from space charge density analysis. They attributed the observed defect level to interstitial Cu- and a V<sub>Cd</sub> pair. Similarly, Al-Bassam et al. reported an additional DLE in Zn0.3Cd0.7 Se crystal at (1.1-1.15) eV located

in the middle of the bandgap. It was associated with an electron trap (Al-Bassam & Elani, 2013).

Photoluminescence spectroscopy (PL) measurements offer the opportunity to study defect states in a semiconductor material. PL emissions at longer wavelengths are usually associated with defect centers. An electron excited from the ground states becomes localized at mid-gap impurity centers. The excess energy is non-radiatively transmitted in the form of heat (vibrations) to the crystal lattice after the time ( $t_n$ ); the process can lead to the generation of lattice defects. The low-energy component of the spectra consists of various emissions from structural defects.

In this work, we aim to reveal for the first time the defect characteristics of milled  $Cd_{0.3}Zn_{0.7}Se$  nanocrystals at (625 nm to 900)nm spectral window using photoluminescence spectroscopy analysis.

#### MATERIALS AND METHOD

The experimental procedure used to synthesize  $Cd_{0.3}Zn_{0.7}Se$  was similar to that reported earlier by (Muh'd et al., 2015). Three elemental constituents (Cd, Zn, and Se) purchased from (Alfa Aesar of 99.9%) purity were sealed in a grinding jar made of stainless steel. The powder was thoroughly ground for 5, 10, and 20 h intervals using a PM100 (Retsch) at 300 revolutions per minute. One gram of the milled powder was dissolved in 5 ml of acetone for photoluminescence measurements using a Perkin-Elmer (LS-55) luminescence spectrophotometer located at Physics Department University Putra Malaysia.

**RESULTS AND DISCUSSION** 



Figure 1 (a–c). Temporal evolutions of PL spectra at room temperature of Cd<sub>0.3</sub>Zn<sub>0.7</sub> Se nanocrystal for a) 5, b) 10, and c) 20 h. The multi-peak Gaussian fitted to the PL spectra consists of three prominent emission bands evolving over the milling time. The S with an arrow indicates a dislocation related band.

The PL spectrum of the 5 h milled nanocrystals powders is shown in Fig. 1(a) has six emission peaks labeled D1 to D6 (from high photon energy to low photon energy of the spectrum). The spectra of the milled samples were fitted with multi-Gaussian peaks; the fitting parameters are summarized in Table 1. In line with the Gaussian fitting results, the photon energies of these peaks related to D1 and D2 are 1.804 and 1.694 eV, respectively. Peaks D3 and D4 are of the same origin with a photon energy of 1.540 eV. Similarly, peaks D5 and D6 have a photon energy of 1.454 eV. The PL spectral bandwidths of peaks D1–D6 are 16, 70, 24, 73, 41, and 69 nm, respectively.



As shown in Fig.1 (b), the fitted PL peak energies of the 10 h milled sample are represented by P1, P2, P3, P4, P5, and P6 with corresponding photon energies of 1.811, 1.738, 1.676, 1.544, 1.486, and 1.450 eV and PL spectral bandwidths of 23, 16, 96, 26, 60, and 34 nm, respectively.



Wavelength (nm)

In Fig.1 (C), the 20 h milled sample fitted photon energies are W1, W2, W3, W4, and W5 with corresponding photon energies of 1.910, 1.804, 1.539, 1.454, and 1.382 eV and spectral bandwidths of 28, 18, 27, 51, and 12 nm, respective

Table 1. Fitting parameters for	(5-20) h milled Cd <sub>0.3</sub> Zn <sub>0.7</sub> Se powders	•
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Peak No. 5h milled	Peak Assignment	Center Maximum(nm)	Maximum Height	FWHM	Peak energy(eV)
Peak 1	D1	687	4.14	16.47	1.80
Peak 2	D2	733	4.39	69.60	1.69
Peak 3	D3	805	5.68	23.74	1.54
Peak 4	D4	805	6.30	73.86	1.54
Peak 5	D5	852	5.60	40.53	1.45
Peak 6	D6	852	0.71	69.30	1.45
10h milled					
Peak 1	P1	685	7.00	22.60	1.81
Peak 2	P2	714	2.21	15.93	1.74
Peak 3	P3	740	6.56	95.77	1.68
Peak 4	P4	803	10.11	26.41	1.54
Peak 5	Р5	834	6.99	59.82	1.50
Peak 6	Р6	855	4.77	34.15	1.45
20h milled					
Peak1	W1	651	6.36	27.91	1.91
Peak2	W2	687	14.26	18.12	1.80

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Peak3	W3	805	18.81	27.22	1.54
Peak4	W4	853	16.90	50.68	1.45
Peak5	W5	897	3.75	12.34	1.38

In the PL spectra of all three samples, a band related to dislocation donated by (S) at approximately 1.7 eV was observed. The integral intensity of this band exhibits approximately five-fold decrease after 20 h of milling. Dislocation-related emission peaks were observed in mechanically alloyed ZnSe and ZnO (Radoi et al., 2003, 2004). The fitted emission peaks in Fig.1 (a–c) are dominated by deep defect level emission from 650–897 nm (1.38–1.91 eV). Similarly, deep-level emission (DL) was observed in ZnSe, CdSe, and CdZnSe, as shown in Table 2. Such emissions are usually attributed to stacking faults and stoichiometric defects

that include intrinsic point defects (vacancies and interstitials). An important mechanism for luminescence in semiconductors is the recombination between an electron trapped at one defect (donor) and a hole trapped at another defect (acceptor). A large concentration of vacancies and interstitial with different charge states ( $V_{Cd}$ ,  $V_{Zn}$ , and  $V_{Se}$ ) are typically found in the milled samples. These defects generate additional localized energy levels within the bandgap and eventually recombine into donor-acceptor pairs (DAPs).

Table 2. Energy position of defect centers and do	nor-acceptor pairs in ZnSe, CdSe, and Cd	ZnSe.
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Material	Defect	Energy level (nm)	Ref.
ZnSe	(Vzn–Zni)	600	(Rong & Watkins, 1986)
		800	
	VZn	720	(Lee et al., 1980)
		635	(Dunstan et al., 1977)(Philipose et al.,
2007)			
		652	(Radoi et al., 2004)(Allen, 1995)
	(Vzn–Vse)	730	(Makhni & Tkachenko, 2003)
	(V <sub>Zn</sub> –Zn <sub>i</sub> )	652	(Makhni & Tkachenko, 2003)
CdSe	V <sub>Cd</sub>	885	(Norris, 1982)
	Vse	861	(Pal et al., 2001)
		704-800	(Jager-Waldau et al., 1988)
CdZnSe		590–750	

The values of DL emissions reported in Table 2 are consistent with the observed values (Table 1). Thus, the DL emission is ascribed to a DAP with zinc or selenium vacancy and zinc interstitial. Similarly, the observed dominant emission peak (805 nm) coincides with the 800 nm luminescence peak as reported by Rong *et al.* (Rong & Watkins, 1986). The peak arises from localized charge transfer or excitonic transition involving Frenkel pairs ( $V_{Zn}$ -Zn<sub>i</sub>). The peak at 885 nm (1.40 eV) was ascribed to cadmium vacancy (Norris, 1982) and is consistent with peak W5 in Table1.

The emission intensities of W2, W3, and W4 show 3.4, 3.3, and a 3.0-fold increase compared with D1, D3, and D4 peaks, respectively. Besides, D1, D3, and D4 peaks show a considerable increase in FWHM values, as shown in Table 2. An increase in PL intensity with decreasing crystallite size (increased milling time) is unpredicted because the concentration of surface state is expected to increase with decreasing crystallite size. Similar observations were already reported [12–14]. The luminescence efficiency  $\eta$  in the nanocrystal can be expressed as $\eta = (1 + \beta D^2)^{-1}$ , where  $\beta$  is related to the ratio of radiative and non-radiative decay time  $\binom{\iota_R}{\tau_{NR}}$  and D is the volume of the nanocrystal. Thus, the luminescence efficiency in nanocrystal increases with decreasing crystallite size. Also, PL intensity is linked to improvement in the surface quality of crystallite upon growth mechanism.

The emission in this spectral window is typically associated with deep-level states arising from stacking faults as well as intrinsic point defects such as vacancies and interstitials. In the parent semiconductor ZnSe, the intrinsic point defects reported are the Zn vacancy  $(V_{Zn})$ , Zn interstitial  $(Zn_i)$ , Se vacancies  $(V_{Se})$ , and Se interstitials (Sei) (see Table 2). However, abundant Vse is not expected because of the large formation energy of Vse and the low melting point of Se, resulting in the rapid disappearance of selenium peaks. Furthermore, the observed intrinsic defects are characterized by different charge states and therefore will produce different localized energy levels in the bandgap of the semiconductor. For instance,  $V_{Zn}$  is a double acceptor level with possible charge states of VZn<sup>-</sup> and VZn<sup>2-</sup>. Thus, Vse vacancies and Zni produce states close to the conduction band and will give rise to acceptor states. On the other hand, the produced  $V_{Zn}$  and Se<sub>i</sub> are donor states near the valance band. The PL emissions of The number of possible intrinsic defects and DAP transitions can give rise to complex PL spectra with a multitude of peaks. The PL spectra for our samples show a minimum of 6 peaks within the measured spectral window. Interestingly, for the samples milled at 5, 10, and 20h, we identify three PL peaks at 1.8, 1.54, and 1.45 eV; the position of the peaks remains constant with milling time. This result is unexpected because the average crystallite diameters reduced with milling time; therefore, the samples milled for longer durations have wider band gaps. A possible explanation is that these PL peaks are associated with electronic transitions from defect levels to the bulk bands. The acceptor defect levels are fixed to the valance band. On the other hand, donor levels have the same features comparable to states in the conduction band that to increase by the energy of the conduction band shift (Oba et al., 2001). Despite the change of band gaps with milling time, the formation energies of acceptor and donor defects remain constant to the conduction and valance band edge. The intensity of these PL peaks increases with milling time because longer milling time corresponds to a higher density of intrinsic defects. A plot of the three PL peaks versus milling time is shown in Fig.2 below.



Figure 2. Dependence of PL energy position of different DAP pairs on milling time.

In addition to the three PL peaks, we observe other broadband peaks with a strong dependence on increased milling time (reduced particle diameter). These peaks have the characteristics of the DAP transition. As broadband emission lines, DAP transition energies are dependent on the spatial separation and transition probability of the pairs involved in the recombination process that leads to PL emission. Thus, the spatial distribution of point defects in the nanocrystals can give an enormous number of closely spaced PL lines, leading to a broad DAP line. Additionally, the peak position of DAP lines shifts as the bandgap changes. This phenomenon will lead to a change in the relative energy of the acceptor and donor pairs. The PL energy is given by Equation (1).

$$E_{PL} \cong E_{Zni} \left( \frac{+}{2} + \right) - E_{V_{Zn}} (2 - / -) + \frac{e^2}{\varepsilon r} + E_{relax}$$
(1)

The first terms in the equation represent the energy positions of the donor and acceptor states. On the other hand, the third term donates Coulomb energy, which decreases with r; this phenomenon is consistent with the behavior of other peaks in which the energy position of the peaks shifts toward the longer wavelength.

CONCLUSION

In summary, defect characterization of mechanically alloyed CdZnSe nanocrystals was conducted using photoluminescence spectroscopy at room temperature. Two prominent emission peaks (1.4, 1.8, and 1.54) eV were detected in the PL spectrum of 650-900 nm. The peaks remain fixed at a particular energy position without a shift to a higher energy level despite a decrease in particle size of the nanocrystals. We argued that the observed peaks could be due to vacancies and interstitial defects with different charge states ( $V_{Cd}$ ,  $V_{Zn}$ , and  $V_{Se}$ ) found in the milled samples.

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